

A posteriori local subcell correction of high-order DG schemes on unstructured grids

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- 1 Introduction
- 2 DG as a subcell Finite Volume
- 3 *A posteriori* subcell correction
- 4 Numerical results
- 5 Conclusion

Scalar conservation law

- $\partial_t u(\mathbf{x}, t) + \nabla_x \cdot \mathbf{F}(u(\mathbf{x}, t)) = 0, \quad (\mathbf{x}, t) \in \omega \times [0, T]$
- $u(\mathbf{x}, 0) = u_0(\mathbf{x}), \quad \mathbf{x} \in \omega$

$(k + 1)^{\text{th}}$ order semi-discretization

- $\{\omega_c\}_c$ a partition of ω , such that $\omega = \bigcup_c \omega_c$
- $u_h(\mathbf{x}, t)$ the numerical solution, such that $u_h|_{\omega_c} = u_h^c \in \mathbb{P}^k(\omega_c)$

$$u_h^c(\mathbf{x}, t) = \sum_{m=1}^{N_k} u_m^c(t) \sigma_m^c(\mathbf{x})$$

- $\{\sigma_m^c\}_{m=1, \dots, N_k}$ a basis of $\mathbb{P}^k(\omega_c)$, with $N_k = \frac{(k+1)(k+2)}{2}$ in 2D.

Local variational formulation on ω_c

- $\int_{\omega_c} \frac{\partial u_h^c}{\partial t} \psi \, dV = \int_{\omega_c} \mathbf{F}(u_h^c) \cdot \nabla_x \psi \, dV - \int_{\partial \omega_c} \psi \mathcal{F}_n \, dS, \quad \forall \psi \in \mathbb{P}^k(\omega_c)$
- $\mathcal{F}_n = \mathcal{F}(u_h^c, u_h^v, \mathbf{n})$ numerical flux

Numerical example: solid body rotation

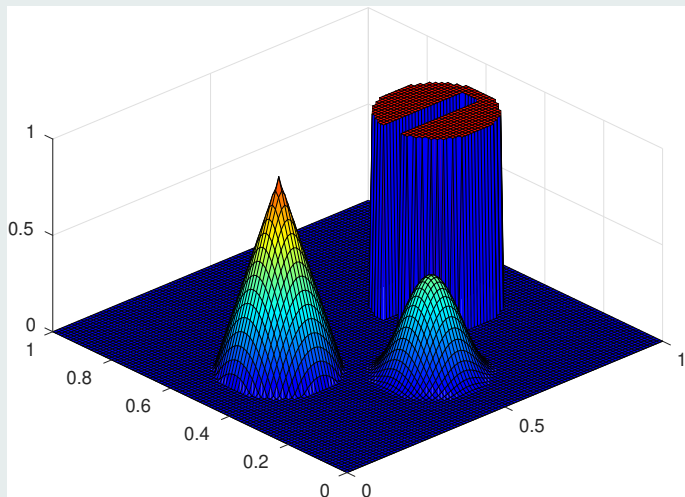
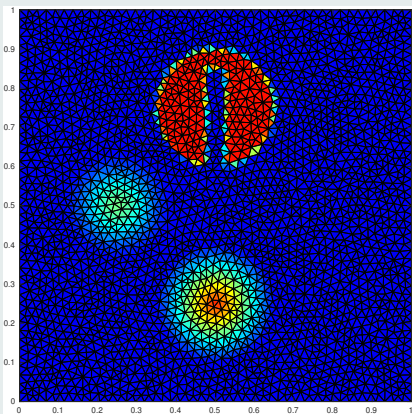
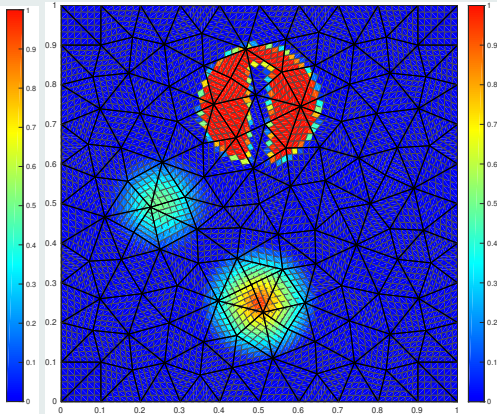


Figure : Rotation of composite signal: initial solution

Roughly constant number of degrees of freedom



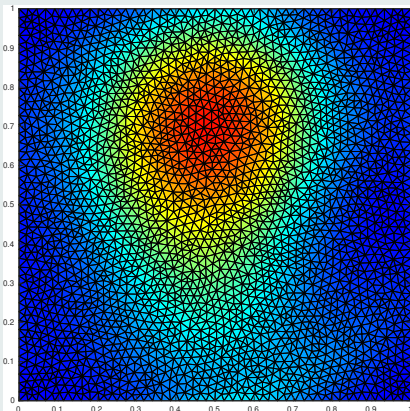
(a) 1st order on 5154 cells



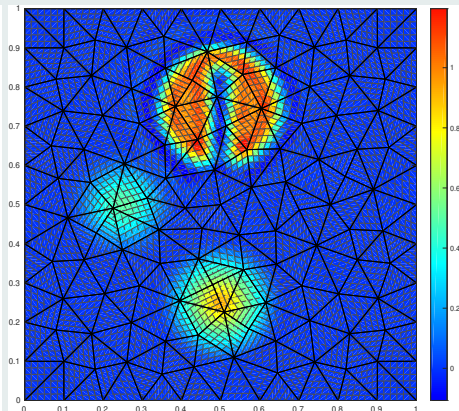
(b) 6th order on 242 cells (5082 DoF)

Figure : Rotation of composite signal: initial solution

Subcell resolution of DG scheme



(c) 1st order on 5154 cells



(d) 6th order on 242 cells (5082 DoF)

Figure : Rotation of composite signal after one period: subcells mean value

Subcell resolution of DG scheme

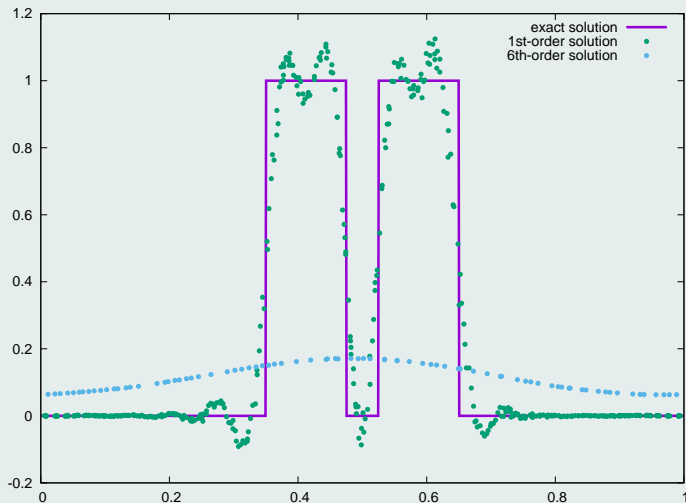


Figure : Rotation of composite signal after one period: profiles for $y = 0.75$

Admissible numerical solution

- Maximum principle / positivity preserving
- Prevent the code from crashing (for instance avoiding NaN)
- **Ensure the conservation of the scheme**

Spurious oscillations

- Discrete maximum principle
- Relaxing condition for smooth extrema

Accuracy

- Retain as much as possible the subcell resolution of the DG scheme
- Minimize the number of subcell solutions to recompute

Modify locally, at the subcell level, the numerical solution without impacting the solution elsewhere in the cell



F. VILAR, *A Posteriori Correction of High-Order DG Scheme through Subcell Finite Volume Formulation and Flux Reconstruction*. JCP, 2018.

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DG as a subcell Finite Volume

- Rewrite DG scheme as a FV-like scheme on a subgrid

Cell subdivision into N_k subcells

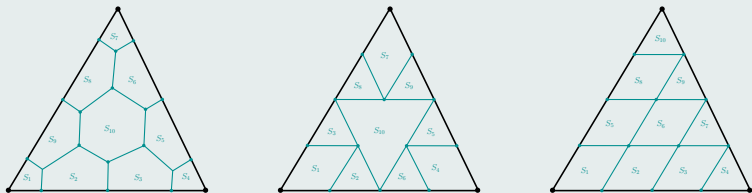


Figure : Examples of subdivision for a \mathbb{P}^3 DG scheme on a triangular cell

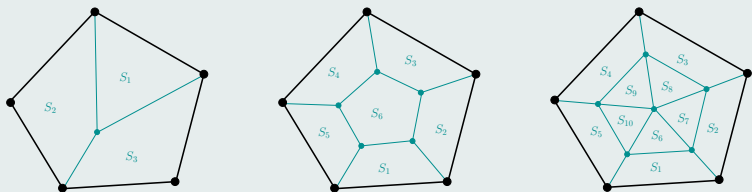


Figure : Examples of subdivision for a polygonal cell from \mathbb{P}^1 up to \mathbb{P}^3

DG schemes through residuals

$$\bullet \sum_{m=1}^{N_k} \frac{d u_m^c}{dt} \int_{\omega_c} \sigma_m \sigma_p dV = \int_{\omega_c} \mathbf{F}(u_h^c) \cdot \nabla_x \sigma_p dV - \int_{\partial\omega_c} \sigma_p \mathcal{F}_n dS, \quad \forall p \in \llbracket 1, N_k \rrbracket$$

$$\implies \boxed{M_c \frac{d U_c}{dt} = \Phi_c}$$

- $(U_c)_m = u_m^c$ Solution moments
- $(M_c)_{mp} = \int_{\omega_c} \sigma_m \sigma_p dV$ Mass matrix
- $(\Phi_c)_m = \int_{\omega_c} \mathbf{F}(u_h^c) \cdot \nabla_x \sigma_m dV - \int_{\partial\omega_c} \sigma_m \mathcal{F}_n dS$ DG residuals

Subdivision and definition

- ω_c is subdivided into N_k subcells S_m^c
- Let us define $\bar{\psi}_m^c = \frac{1}{|S_m^c|} \int_{S_m^c} \psi dV$ the subcell mean value

Submean values

$$\bullet \frac{d\bar{u}_m^c}{dt} = \frac{1}{|S_m^c|} \sum_{q=1}^{N_k} \frac{d u_q^c}{dt} \int_{S_m^c} \sigma_q dV$$

 \Rightarrow

$$\frac{d\bar{U}_c}{dt} = P_c \frac{dU_c}{dt}$$

$$\bullet (\bar{U}_c)_m = \bar{u}_m^c$$

Submean values

$$\bullet (P_c)_{mp} = \frac{1}{|S_m^c|} \int_{S_m^c} \sigma_p dV$$

Projection matrix

 \Rightarrow

$$\frac{d\bar{U}_c}{dt} = P_c M_c^{-1} \Phi_c$$

Admissibility of the cell sub-partition into subcells

- P_c has to be non-singular

Subcell Finite Volume: reconstructed fluxes

- Let us introduce the **reconstructed fluxes** such that

$$\frac{d\bar{u}_m^c}{dt} = -\frac{1}{|S_m^c|} \int_{\partial S_m^c} \widehat{F}_n dS$$

- Let \mathcal{V}_m^c be the set of face neighboring subcells of S_m^c

$$\frac{d\bar{u}_m^c}{dt} = -\frac{1}{|S_m^c|} \sum_{S_p^v \in \mathcal{V}_m^c} \int_{f_{mp}^c} \widehat{F}_n dS$$

- We impose that on the boundary of cell ω_c

$$\widehat{F}_n|_{\partial\omega_c} = \mathcal{F}_n$$

- Then, if $\widetilde{\mathcal{V}}_m^c$ stands for the set of face neighboring subcells inside ω_c

$$\frac{d\bar{u}_m^c}{dt} = -\frac{1}{|S_m^c|} \left(\sum_{S_p^v \in \widetilde{\mathcal{V}}_m^c} \int_{f_{mp}^c} \widehat{F}_n dS + \int_{\partial S_m^c \cap \partial\omega_c} \mathcal{F}_n dS \right)$$

Subcell Finite Volume: reconstructed fluxes

- Taking two subcells S_m^c and S_p^v , the orientation face function ε_{mp}^c writes

$$\varepsilon_{mp}^c = \begin{cases} 1 & \text{if face } f_{mp}^c \text{ is direct or if } f_{mp}^c \subset \partial\omega_c, \\ -1 & \text{if face } f_{mp}^c \text{ is indirect,} \\ 0 & \text{if } S_p^v \notin \mathcal{V}_m^c. \end{cases}$$

- $\int_{f_{mp}^c} \widehat{F}_n \, dS = \varepsilon_{mp}^c \widehat{F}_{mp}$ face integrated reconstructed flux

$$\frac{d\bar{U}_m^c}{dt} = -\frac{1}{|S_m^c|} \left(\sum_{S_p^v \in \widetilde{\mathcal{V}}_m^c} \varepsilon_{mp}^c \widehat{F}_{mp} + \int_{\partial S_m^c \cap \partial\omega_c} \mathcal{F}_n \, dS \right)$$

- $(B_c)_m = \int_{\partial S_m^c \cap \partial\omega_c} \mathcal{F}_n \, dS$ Cell boundary contribution
- $(A_c)_{mp} = \varepsilon_{mp}^c$ Adjacency matrix
- $D_c = \text{diag}(|S_1^c|, \dots, |S_{N_k}^c|)$ Subcells volume matrix

Subcell Finite Volume: reconstructed fluxes

- Let \widehat{F}_c be the vector containing all the interior faces reconstructed fluxes
- The subcell mean values governing equations yield the following system

$$-A_c \widehat{F}_c = D_c \frac{d\bar{U}_c}{dt} + B_c$$

Graph Laplacian technique

- $A_c \in \mathcal{M}_{N_k \times N_f^c}$ with N_f^c the number of interior faces
- $A_c^t \mathbf{1} = \mathbf{0}$ where $\mathbf{1} = (1, \dots, 1)^t \in \mathbb{R}^{N_k}$



R. ABGRALL, *Some Remarks about Conservation for Residual Distribution Schemes*. *Methods Appl. Math.*, 18:327-351, 2018.

- Let \mathcal{L}_c^{-1} be the inverse of $L_c = A_c A_c^t$ on the orthogonal of its kernel

$$\mathcal{L}_c^{-1} = (L_c + \lambda \Pi)^{-1} - \frac{1}{\lambda} \Pi$$

$$\forall \lambda \neq 0$$

- $\Pi = \frac{1}{N_k} (\mathbf{1} \otimes \mathbf{1}) \in \mathcal{M}_{N_k}$

Graph Laplacian technique

- Finally, we obtain the following definition of the reconstructed fluxes

$$\widehat{F}_c = -A_c^t \mathcal{L}_c^{-1} \left(D_c P_c M_c^{-1} \Phi_c + B_c \right)$$

remark

- The only terms depending on the time are Φ_c and B_c

Back to the DG scheme

- The polynomial solution is defined through reconstructed fluxes as follows

$$\frac{dU_c}{dt} = -P_c^{-1} D_c^{-1} \left(A_c \widehat{F}_c + B_c \right)$$

Question

- Is the reconstructed flux \widehat{F}_c close to the interior flux $F(u_h^c)$?

Local variational formulation

- $$\int_{\omega_c} \frac{\partial u_h^c}{\partial t} \psi \, dV = \int_{\omega_c} \mathbf{F}(u_h^c) \cdot \nabla_x \psi \, dV - \int_{\partial\omega_c} \psi \mathcal{F}_n \, dS, \quad \forall \psi \in \mathbb{P}^k(\omega_c)$$
- Substitute $\mathbf{F}(u_h^c)$ with $\mathbf{F}_h^c \in (\mathbb{P}^{k+1}(\omega_c))^2$ (collocated or L_2 projection)
- $$\int_{\omega_c} \frac{\partial u_h^c}{\partial t} \psi \, dV = - \int_{\omega_c} \psi \nabla_x \cdot \mathbf{F}_h^c \, dV + \int_{\partial\omega_c} \psi (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS, \quad \forall \psi \in \mathbb{P}^k(\omega_c)$$

Subresolution basis functions

- Let us introduce the N_k basis functions $\{\phi_m\}_m$ such that $\forall \psi \in \mathbb{P}^k(\omega_c)$

$$\int_{\omega_c} \phi_m \psi \, dV = \int_{S_m^c} \psi \, dV, \quad \forall m = 1, \dots, N_k,$$

- $$\sum_{m=1}^{N_k} \phi_m(\mathbf{x}) = 1$$

These particular functions can be seen as the L_2 projection of the indicator functions $\mathbb{1}_m(\mathbf{x})$ onto $\mathbb{P}^k(\omega_c)$

Subcell Finite Volume scheme

$$\bullet \int_{\omega_c} \frac{\partial u_h^c}{\partial t} \phi_m \, dV = - \int_{\omega_c} \phi_m \nabla_x \cdot \mathbf{F}_h^c \, dV + \int_{\partial\omega_c} \phi_m (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS$$

$$\bullet |S_m^c| \frac{d\bar{u}_m^c}{dt} = - \int_{S_m^c} \nabla_x \cdot \mathbf{F}_h^c \, dV + \int_{\partial\omega_c} \phi_m (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS$$

$$\bullet \frac{d\bar{u}_m^c}{dt} = - \frac{1}{|S_m^c|} \left(\int_{\partial S_m^c} \mathbf{F}_h^c \cdot \mathbf{n} \, dS - \int_{\partial\omega_c} \phi_m (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS \right)$$

$$\bullet \frac{d\bar{u}_m^c}{dt} = - \frac{1}{|S_m^c|} \int_{\partial S_m^c} \widehat{F}_n \, dS$$

Subcell Finite Volume

Reconstructed Fluxes

- Finally, we get that

$$\int_{\partial S_m^c} \widehat{F}_n \, dS = \int_{\partial S_m^c} \mathbf{F}_h^c \cdot \mathbf{n} \, dS - \int_{\partial\omega_c} \phi_m (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS$$

Reconstructed fluxes

- As we impose that $\widehat{F}_n|_{\partial\omega_c} = \mathcal{F}_n$, this last expression rewrites

$$\int_{\partial S_m^c \setminus \partial\omega_c} \widehat{F}_n \, dS = \int_{\partial S_m^c \setminus \partial\omega_c} \mathbf{F}_h^c \cdot \mathbf{n} \, dS - \int_{\partial\omega_c} \widetilde{\phi}_m (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS$$

- $$\widetilde{\phi}_m = \begin{cases} \phi_m & \text{if } \mathbf{x} \in \partial\omega_c \setminus \partial S_m^c \\ \phi_m - 1 & \text{if } \mathbf{x} \in \partial\omega_c \cap \partial S_m^c \end{cases}$$

- $$\int_{f_{mp}^c} \widehat{F}_n \, dS = \varepsilon_{mp}^c \widehat{F}_{mp} \quad \text{and} \quad \int_{f_{mp}^c} \mathbf{F}_h^c \cdot \mathbf{n} \, dS = \varepsilon_{mp}^c F_{mp}$$

- Then, if F_c is the vector containing all the interior faces fluxes, one gets

$$A_c \widehat{F}_c = A_c F_c - G_c$$

- $$(G_c)_m = \int_{\partial\omega_c} \widetilde{\phi}_m (\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n) \, dS$$

Boundary contribution

Reconstructed fluxes through interior fluxes

- Making use of the same graph Laplacian technique, we finally obtain

$$\widehat{F}_c = F_c - A_c^t \mathcal{L}_c^{-1} G_c$$

- We can rewrite this expression as

$$\widehat{F}_c = F_c - \mathcal{G}(\mathbf{F}_h^c \cdot \mathbf{n} - \mathcal{F}_n)$$

where $\mathcal{G}(\cdot)$ is a correction function taking into account the jump between the polynomial flux and the numerical flux on the cell boundary

Remark

- Different choice in the correction term $\mathcal{G}(\cdot)$ leads to different schemes
- For instance, $\mathcal{G}(\cdot) = 0$ leads to the spectral volume scheme of Z.J. Wang

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RKDG scheme

- SSP Runge-Kutta: convex combinations of first-order forward Euler
- For sake of clarity, we focus on forward Euler time stepping

Projection on subcells of RKDG solution

- $u_h^{c,n}(x) = \sum_{m=1}^{N_k} u_m^{c,n} \sigma_m(x)$
- $u_h^{c,n}$ is uniquely defined by its N_k submean values $\bar{u}_m^{c,n}$
- Recalling the definition of the projection matrix $(P_c)_{mp} = \frac{1}{|S_m^c|} \int_{S_m^c} \sigma_p dV$,

$$\implies P_c \begin{pmatrix} u_1^{c,n} \\ \vdots \\ u_{N_k}^{c,n} \end{pmatrix} = \begin{pmatrix} \bar{u}_1^{c,n} \\ \vdots \\ \bar{u}_{N_k}^{c,n} \end{pmatrix}$$

Set up

- We assume that, for each cell, the $\{\bar{u}_m^{c,n}\}_m$ are admissible
- Compute a candidate solution u_h^{n+1} from u_h^n through uncorrected DG
- For each subcell, check if the submean values $\{\bar{u}_m^{c,n+1}\}_m$ are OK

Physical admissibility detection (PAD)

- Check if $\bar{u}_m^{c,n+1}$ lies in an convex physical admissible set (maximum principle for SCL, positivity of the pressure and density for Euler, ...)
- Check if there is any NaN values

Numerical admissibility detection (NAD)

- Discrete maximum principle DMP on submean values:

$$\min_{v \in \mathcal{V}(S_m^c)} (\bar{u}_v^n) \leq \bar{u}_m^{c,n+1} \leq \max_{v \in \mathcal{V}(S_m^c)} (\bar{u}_v^n)$$

- $\mathcal{V}(S_m^c)$ set of neighboring subcells of S_m^c , including subcell S_m^c
- **This criterion needs to be relaxed to preserve smooth extrema**

Fundamental principle

- On non-admissible subcell boundaries

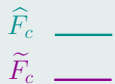
Substitute the reconstructed fluxes by more robust numerical fluxes

- Recompute the non-admissible subcells, and their first neighbors

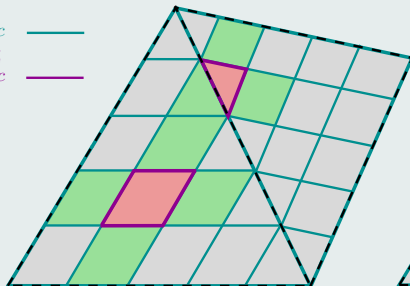
Examples of correction schemes

- **1st-order Finite Volume scheme**
- 2nd-order MUSCL scheme
- (W)ENO methods
- ...

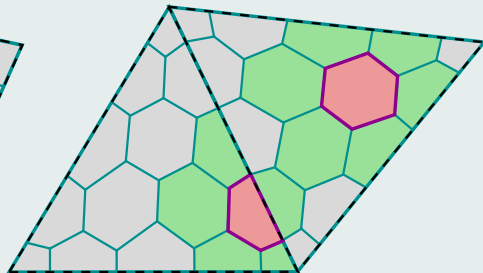
Corrected reconstructed flux

$$\hat{F}_c$$


$$\tilde{F}_c$$



(d) Structured subdivision.



(e) Voronoi-type subdivision.

Figure : Original correction of the DG reconstructed flux

Flowchart

- 1 Compute the uncorrected DG candidate solution $u_h^{c,n+1}$
- 2 Project $u_h^{c,n+1}$ to get the submean values $\bar{u}_m^{c,n+1}$
- 3 Check $\bar{u}_m^{c,n+1}$ through the troubled zone detection plus relaxation
- 4 If $\bar{u}_m^{c,n+1}$ is admissible, go further in time. Otherwise, if S_m^c or $S_p^v \in \mathcal{V}_m^c$ is either marked

$$\widetilde{F}_{mp} = \varepsilon_{mp}^c I_{mp}^c \mathcal{F}(\bar{u}_m^{c,n}, \bar{u}_p^{v,n}, \mathbf{n}_{mp})$$

- 5 Through the corrected reconstructed flux, recompute the submean values for tagged subcells and their first neighbors
- 6 Return to 3

Conclusion

- The correction only affects the DG solution at the subcell scale
- The corrected scheme is conservative at the subcell level
- In practice, few submean values need to be recomputed

Remarks

- For non-linear problems, using very high-order schemes and coarse meshes, the solution may remain a bit oscillatory at the subcell level
- This is why we were previously considering, for $k \geq 3$, that if a subcell is marked as bad then we also mark its first neighboring subcells



F. VILAR, *A Posteriori Correction of High-Order DG Scheme through Subcell Finite Volume Formulation and Flux Reconstruction*. JCP, 2018.

New correction principle

To avoid too much discrepancy between corrected and reconstructed fluxes

- Wider subcell set to be corrected
- Convex combination between 1st-order flux and the reconstructed flux

$$\widetilde{F}_{mp} = \theta_{mp} \varepsilon_{mp}^c I_{mp}^c \mathcal{F}(\bar{u}_m^{c,n}, \bar{u}_p^{v,n}, \mathbf{n}_{mp}) + (1 - \theta_{mp}) \widehat{F}_{mp},$$

where θ_{mp} is a function of the distance to the non-admissible subcell

Corrected reconstructed flux

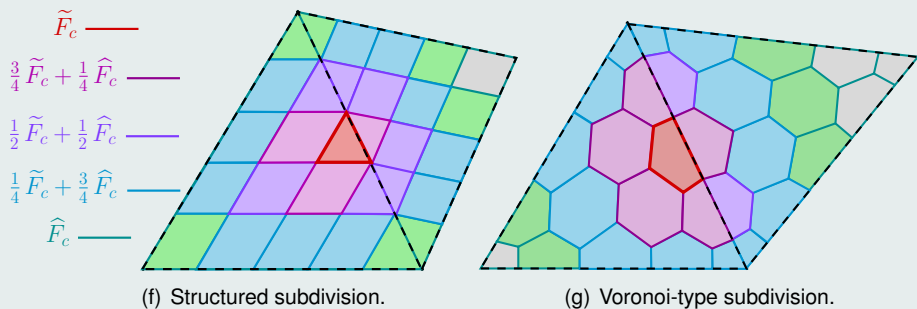
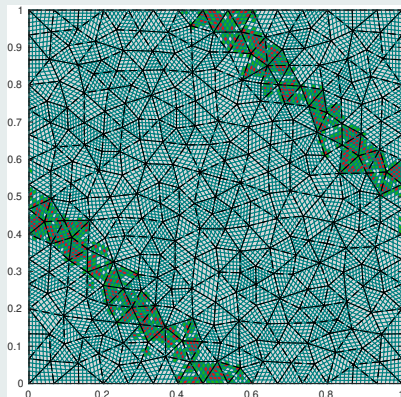


Figure : New correction of the DG reconstructed flux

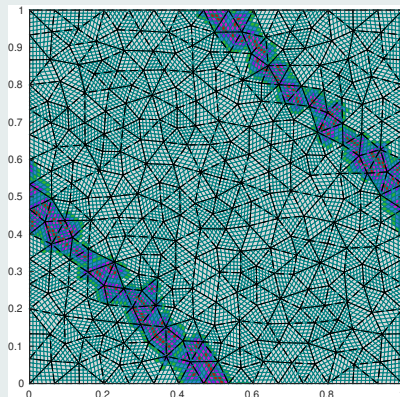
Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$

Figure : Entropic weak solution: apparition of stationary shocks

6th-order scheme on a 576 cells grid

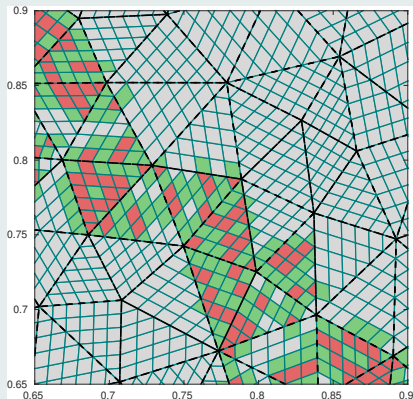


(a) Original correction

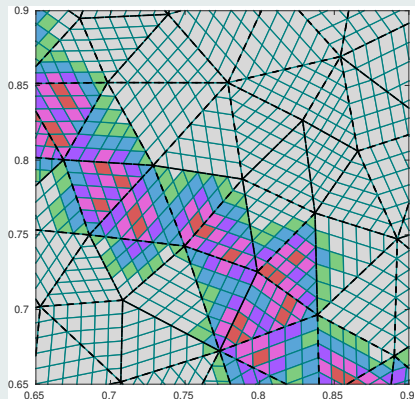


(b) New correction

Figure : Comparison between original and new correction procedure:
corrected subcells

6th-order scheme on a 576 cells grid - zoom in $[0.65, 0.9]^2$ 

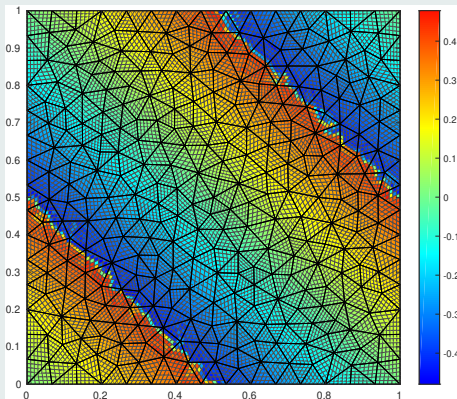
(a) Original correction



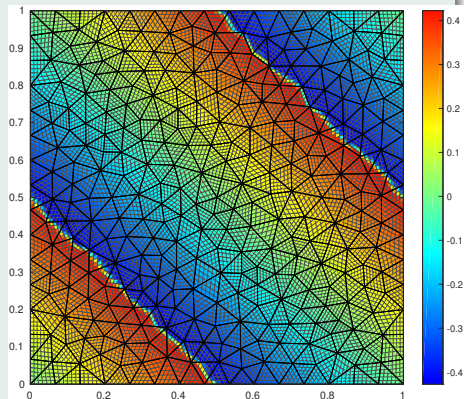
(b) New correction

Figure : Comparison between original and new correction procedure:
corrected subcells

6th-order scheme on a 576 cells grid



(a) Original correction



(b) New correction

Figure : Comparison between original and new correction procedure: subcell mean values

6th-order scheme on a 576 cells grid

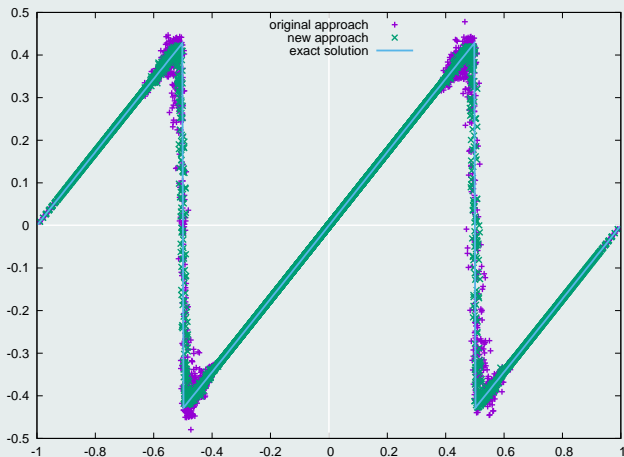


Figure : Comparison between original and new correction procedure: submean values versus $(x + y - 1)$ coordinate

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 - 2D linear problems
 - 2D non-linear problems
- 5 Conclusion

2D Linear advection

- $\partial_t u(\mathbf{x}, t) + \mathbf{A} \cdot \nabla_{\mathbf{x}} u(\mathbf{x}, t) = 0$ with \mathbf{A} transport velocity

Linear advection of a crenel signal

$$\mathbf{A} = (1, 1)^t$$

(a) Subcell mean values

(b) Corrected subcells

Figure : 6th-order APLSC-DG on a 576 cells mesh after one period

Linear advection equation of a crenel signal

$$\mathbf{A} = (1, 1)^t$$

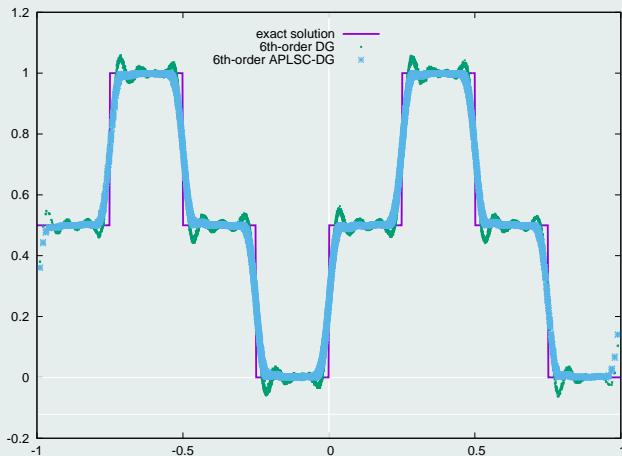
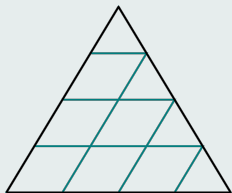
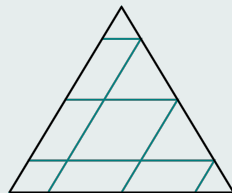


Figure : 6th-order solutions for the crenel advection case on 576 cells: submean values versus $(x + y - 1)$ coordinate

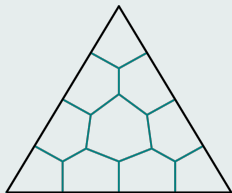
Influence of the subdivision



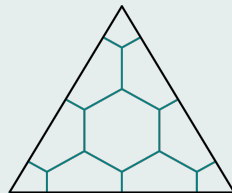
(a) Equidistant boundary points



(b) Gauss-Lobatto boundary points



(c) Equidistant boundary points

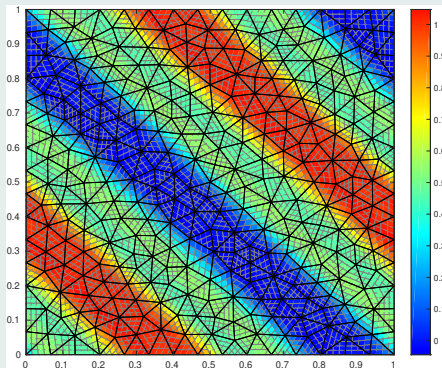


(d) \mathbb{P}^3 Lagrangian mid-points

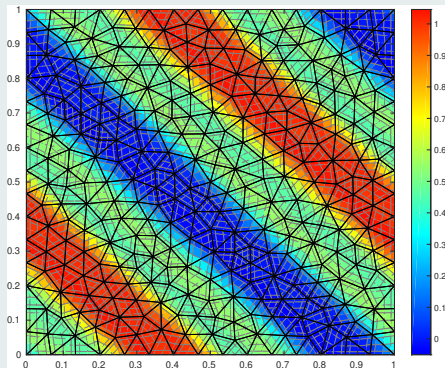
Figure : Examples of subdivisions for a triangular cell and a \mathbb{P}^3 DG scheme

Linear advection equation of a crenel signal

$$\mathbf{A} = (1, 1)^t$$



(a) Uniform structured subdivision

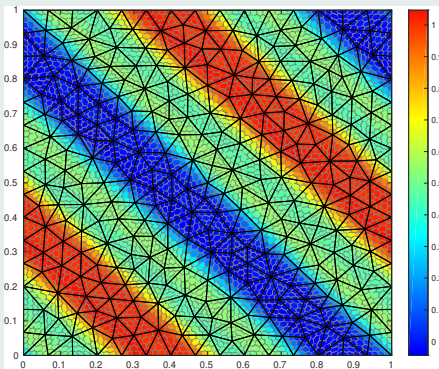


(b) Non-uniform structured subdivision

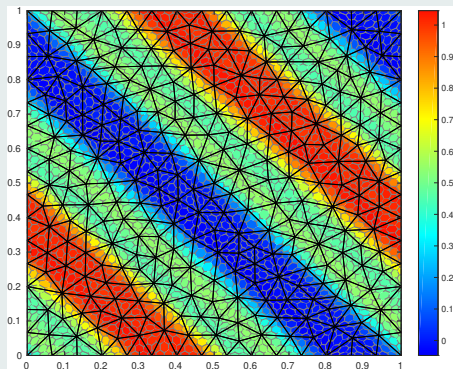
Figure : 4th-order DG solutions for the crenel signal advection on 576 cells after five periods: structured subdivision

Linear advection equation of a crenel signal

$$\mathbf{A} = (1, 1)^t$$



(a) Uniform polygonal subdivision



(b) Non-uniform polygonal subdivision

Figure : 4th-order DG solutions for the crenel signal advection on 576 cells after five periods: polygonal subdivision

Linear advection equation of a crenel signal

$$\mathbf{A} = (1, 1)^t$$

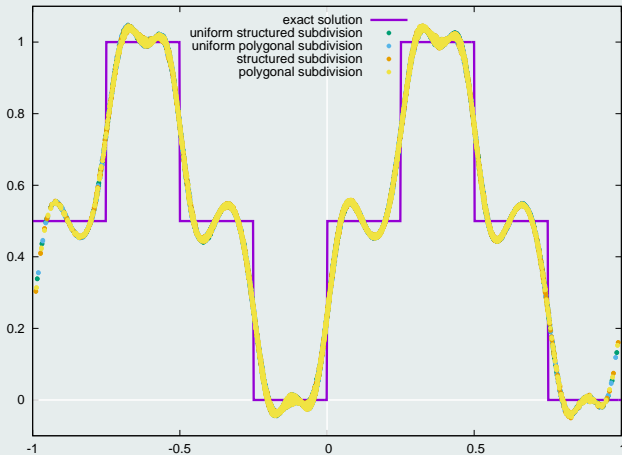
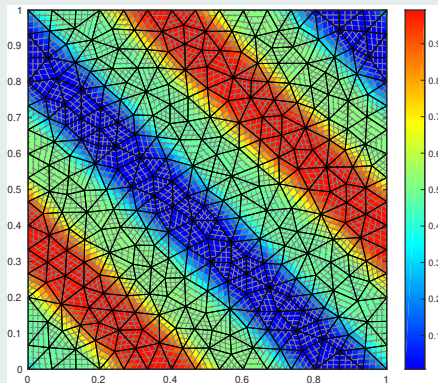


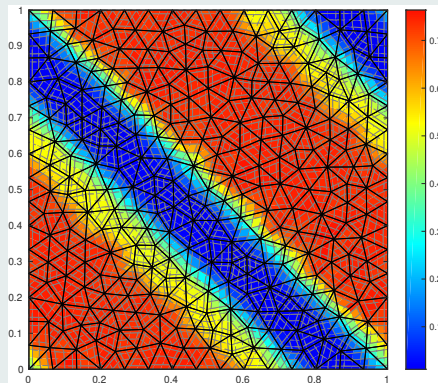
Figure : 4th-order DG solutions for the crenel signal advection on 576 cells using different cell subdivisions: submean values versus $(x + y - 1)$ coordinate

Linear advection equation of a crenel signal

$$\mathbf{A} = (1, 1)^t$$



(a) Uniform structured subdivision

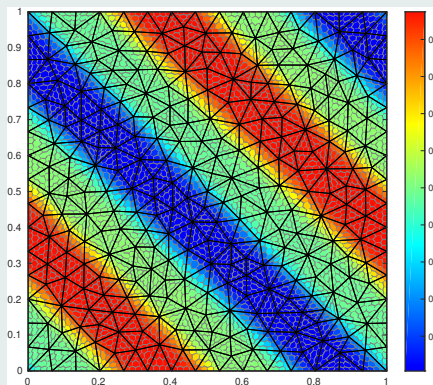


(b) Non-uniform structured subdivision

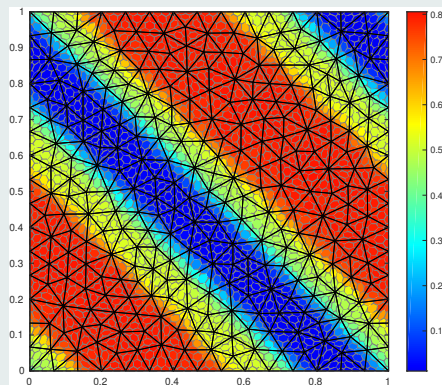
Figure : 4th-order APLSC-DG solutions for the crenel signal advection on 576 cells after five periods: structured subdivision

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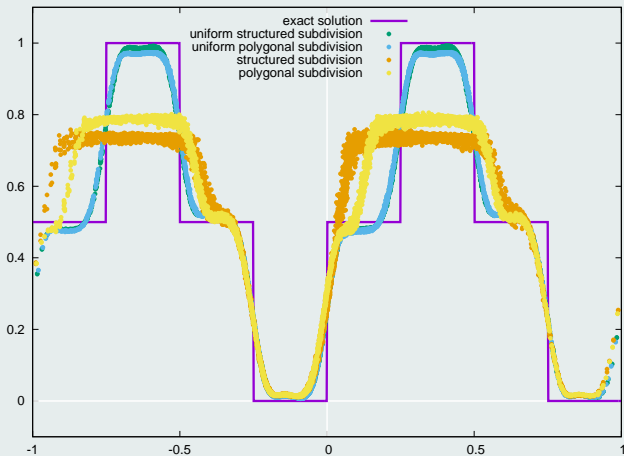


Figure : 4th-order APLSC-DG solutions for the crenel signal advection on 576 cells using different cell subdivisions: submean values versus $(x + y - 1)$ coordinate

2D solid body rotation

- $\partial_t u(\mathbf{x}, t) + \mathbf{A}(\mathbf{x}) \cdot \nabla_{\mathbf{x}} u(\mathbf{x}, t) = 0$ with $\mathbf{A}(\mathbf{x}) = (0.5 - y, x - 0.5)^t$
- $u(\mathbf{x}, 0) = u_0(\mathbf{x})$

Composite signal rotation: 6th-order APLSC-DG on 576 cells

(a) Solution map

(b) Corrected subcells

Rotation of a composite signal after 1 period

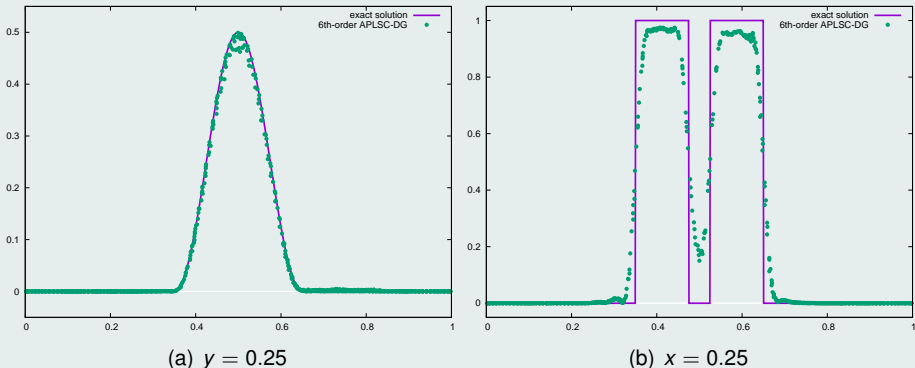
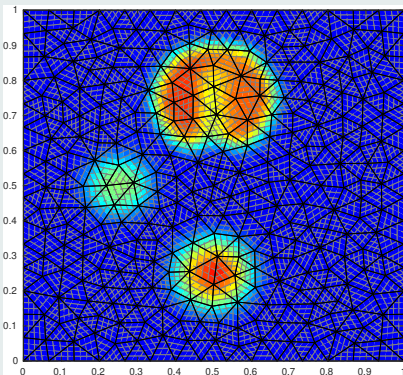
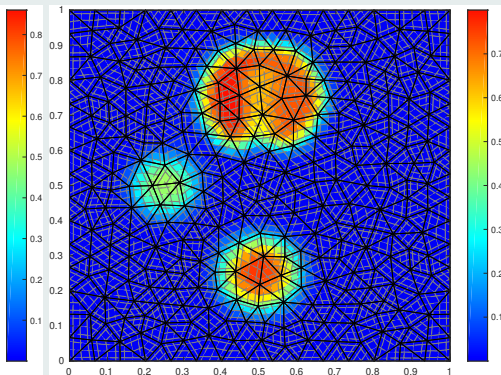


Figure : 6th-order APLSC-DG solution for the rigid rotation case on 576 cells after one full rotation: solution profiles

Rotation of a composite signal after 5 periods



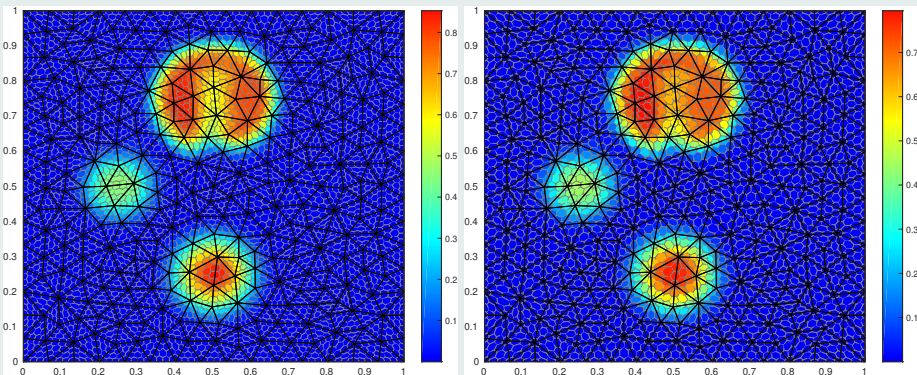
(a) Uniform structured subdivision



(b) Non-uniform structured subdivision

Figure : 4th-order APLSC-DG solutions for the rigid rotation case on 576 cells after five full rotations: structured subdivision

Rotation of a composite signal after 5 periods



(a) Uniform polygonal subdivision

(b) Non-uniform polygonal subdivision

Figure : 4th-order APLSC-DG solutions for the rigid rotation case on 576 cells after five full rotations: polygonal subdivision

Rotation of a composite signal after 5 periods

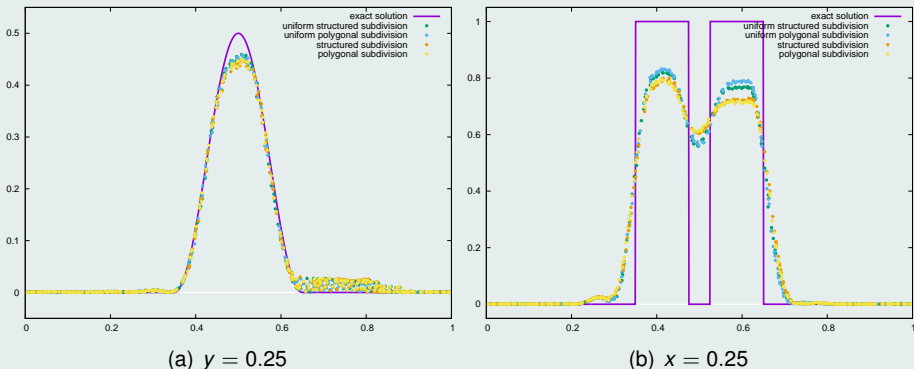


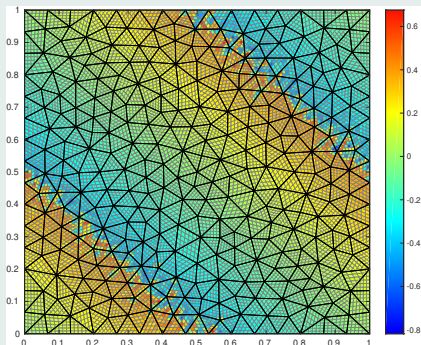
Figure : 4th-order APLSC-DG solutions for the rigid rotation case on 576 cells after five full rotations: solution profiles

- 1 Introduction
- 2 DG as a subcell Finite Volume
- 3 *A posteriori* subcell correction
- 4 Numerical results**
 - 2D linear problems
 - **2D non-linear problems**
- 5 Conclusion

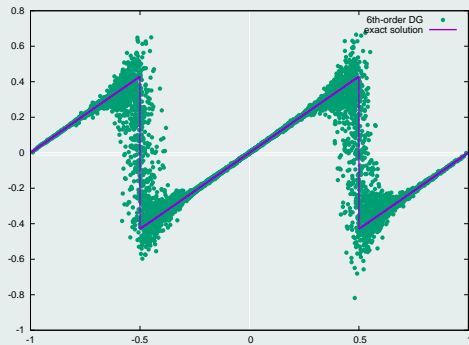
2D non-linear Burgers equation

- $\partial_t u(\mathbf{x}, t) + \nabla_{\mathbf{x}} \cdot \mathbf{F}(u(\mathbf{x}, t)) = 0$ with $\mathbf{F}(u) = \frac{1}{2} (u^2, u^2)^t$
- $u(\mathbf{x}, 0) = u_0(\mathbf{x})$

Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$



(a) Solution map



(b) Solution profile

Figure : 6th-order uncorrected DG on a 576 cells mesh at $t = 0.5$

Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$

(a) Solution map

(b) Corrected subcells

Figure : 6th-order APLSC-DG on a 576 cells mesh at $t = 0.5$

Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$

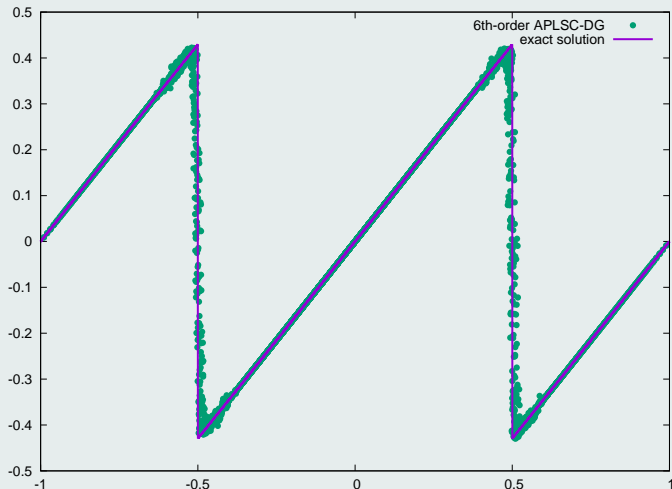
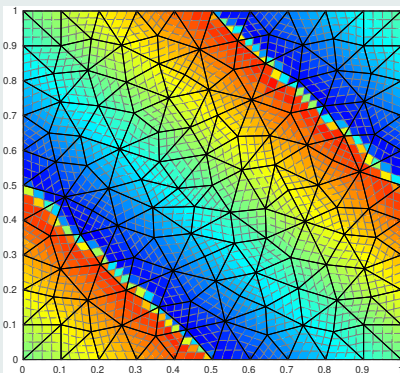
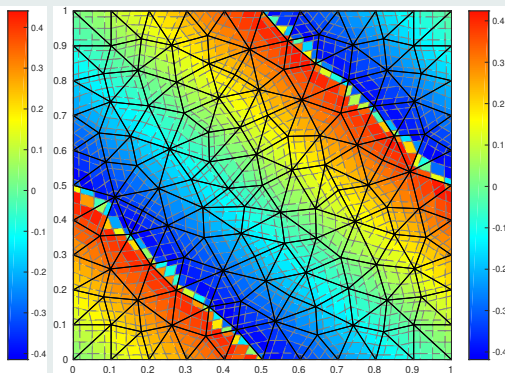


Figure : 6th-order uncorrected DG on a 576 cells mesh at $t = 0.5$: submean values versus $(x + y - 1)$ coordinate

Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$



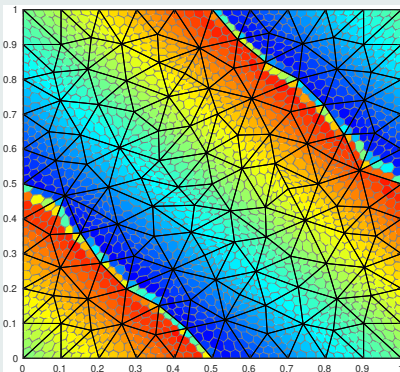
(a) Uniform structured subdivision



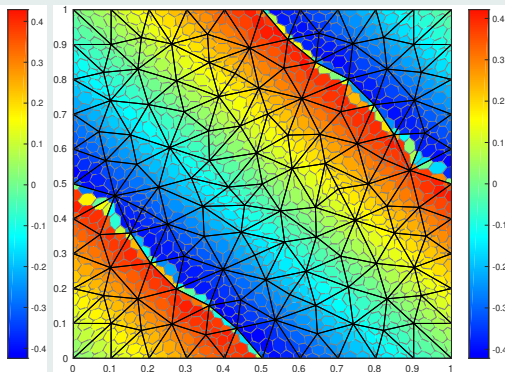
(b) Non-uniform structured subdivision

Figure : 4th-order APLSC-DG solutions for 2D Burgers equation on 242 cells at $t = 0.5$: structured subdivision

Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$



(a) Uniform polygonal subdivision



(b) Non-uniform polygonal subdivision

Figure : 4th-order APLSC-DG solutions for 2D Burgers equation on 242 cells at $t = 0.5$: polygonal subdivision

Burgers equation with $u_0(x, y) = \sin(2\pi(x + y))$

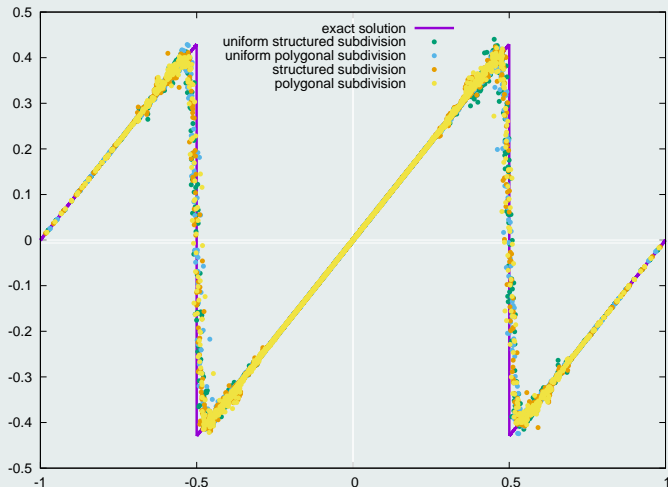


Figure : 4th-order APLSC-DG solutions for 2D Burgers equation on 242 cells at $t = 0.5$: submean values versus $(x + y - 1)$ coordinate

2D Kurganov, Petrova, Popov (KPP) non-convex flux equation

- $\partial_t u(\mathbf{x}, t) + \nabla_{\mathbf{x}} \cdot \mathbf{F}(u(\mathbf{x}, t)) = 0$ with $\mathbf{F}(u) = (\sin u, \cos u)^t$
- $u(\mathbf{x}, 0) = u_0(\mathbf{x})$

KPP non-convex flux problem

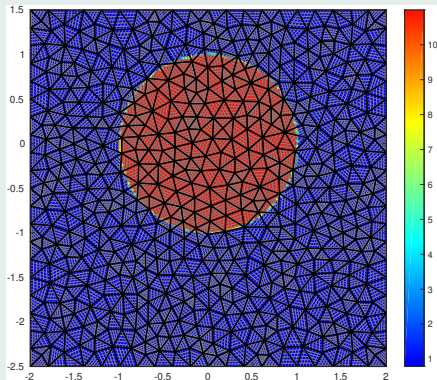
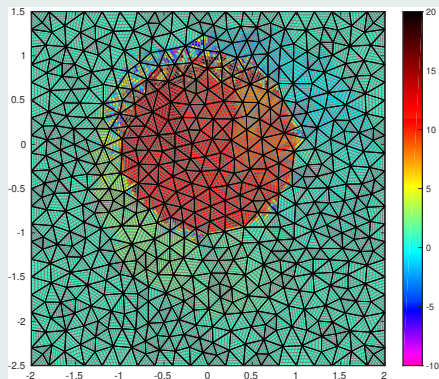
(a) At time $t = 0$ (b) At time $t = 1$

Figure : 6th-order uncorrected DG solution on a 1054 cells mesh

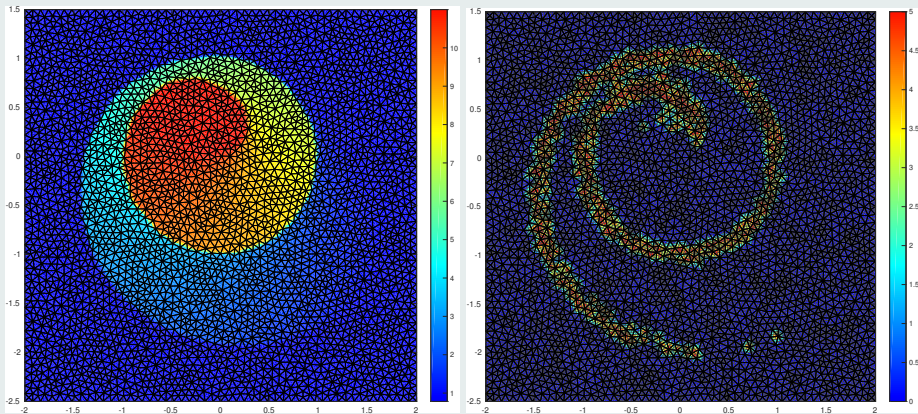
KPP non-convex flux problem

(a) Solution map

(b) Corrected subcells

Figure : 6th-order APLSC-DG solution on a 1054 cells mesh

KPP non-convex flux problem

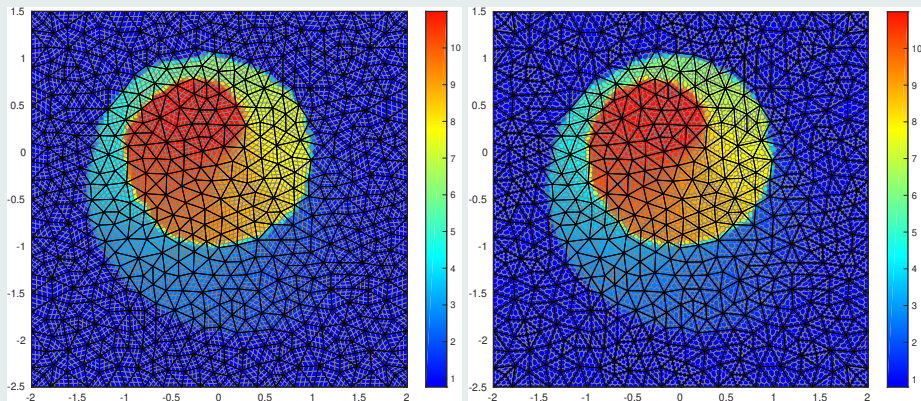


(a) Solution map

(b) Corrected subcells

Figure : 3rd-order APLSC-DG solution on a 6670 cells mesh

KPP non-convex flux problem



(a) Structured subdivision

(b) Voronoi-type subdivision

Figure : 4th-order APLSC-DG solution on a 1054 cells mesh: subdivision comparison

2D non-linear Euler compressible gas dynamics equations

- $\partial_t \mathbf{V} + \nabla_x \cdot \mathbf{F}(\mathbf{V}) = \mathbf{0}$

- $\mathbf{V} = \begin{pmatrix} \rho \\ \mathbf{q} \\ E \end{pmatrix}$

conservative variables

- $\mathbf{F}(\mathbf{V}) = \begin{pmatrix} \mathbf{q} \\ \frac{\mathbf{q} \otimes \mathbf{q}}{\rho} + p I_d \\ (E + p) \frac{\mathbf{q}}{\rho} \end{pmatrix}$

flux function

- $p := p(\mathbf{V}) = (\gamma - 1) \left(E - \frac{1}{2} \frac{\|\mathbf{q}\|^2}{\rho} \right)$

equation of state

APLSC-DG scheme property

- Positivity of the density and internal energy, at the subcell scale

Sod shock tube problem in cylindrical geometry

(a) Density map

(b) Corrected subcells

Figure : 6th-order APLSC-DG solution on a 230 cells mesh

Sod shock tube problem in cylindrical geometry

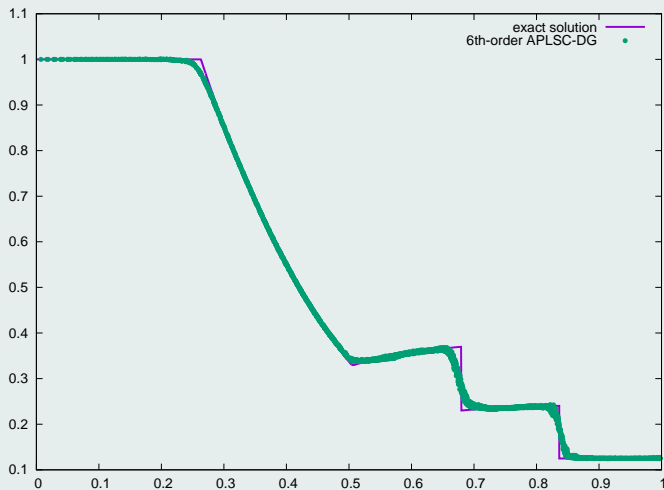
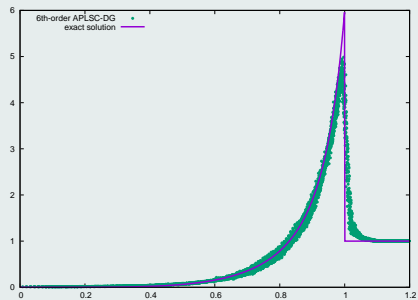


Figure : 6th-order APLSC-DG solution on a 230 cells: density submean values

Sedov point blast problem in cylindrical geometry



(b) Density profile

(a) Energy map

Figure : 6th-order APLSC-DG on a 271 cells mesh at $t = 1$

A Mach 3 wind tunnel with a step

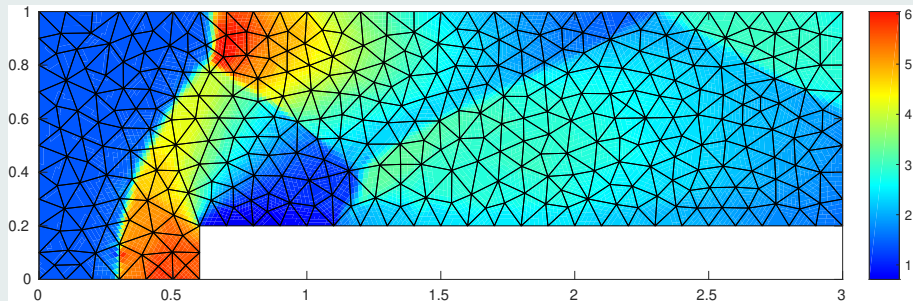


Figure : 6th-order APLSC-DG solution for the facing step problem on 680 cells at $t = 4$: submean density map

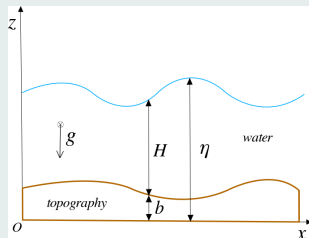
2D non-linear shallow water equations - prebalanced formulation

- $\partial_t \mathbf{V} + \nabla_x \cdot \mathbf{F}(\mathbf{V}, b) = \mathbf{B}(\mathbf{V}, \nabla_x b)$

- $\mathbf{V} = \begin{pmatrix} \eta \\ \mathbf{q} \end{pmatrix}$ conservative variables

- $\mathbf{B}(\mathbf{V}, \partial_x b) = \begin{pmatrix} 0 \\ -g\eta \nabla_x b \end{pmatrix}$ source term

- $\mathbf{F}(\mathbf{V}, b) = \begin{pmatrix} \mathbf{q} \\ \frac{\mathbf{q} \otimes \mathbf{q}}{\eta - b} + \frac{1}{2}g(\eta^2 - 2\eta b) I_d \end{pmatrix}$ flux function



APLSC-DG scheme properties

- Positivity-preservation of the water height $H = \eta - b$, at the subcell scale
- Well-balancing property, at the subcell scale

Well-balancing property

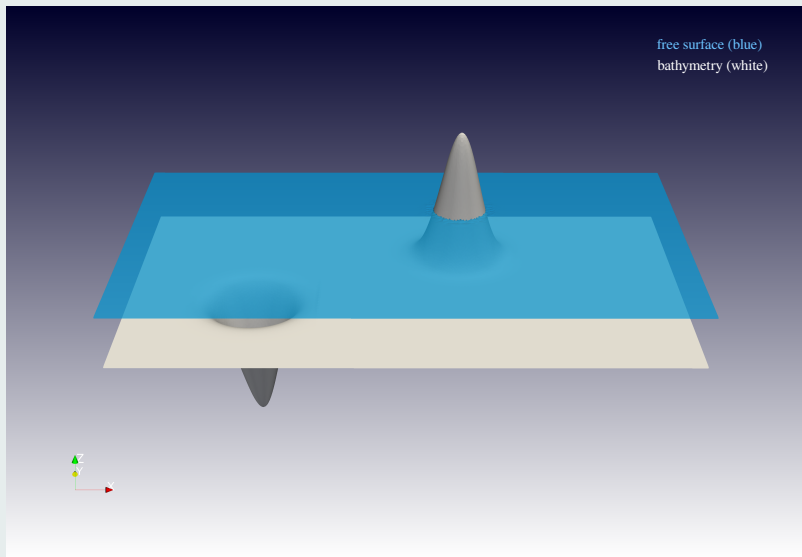
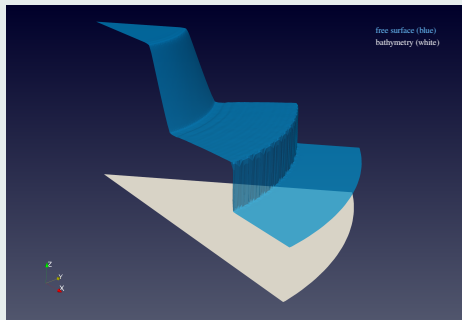
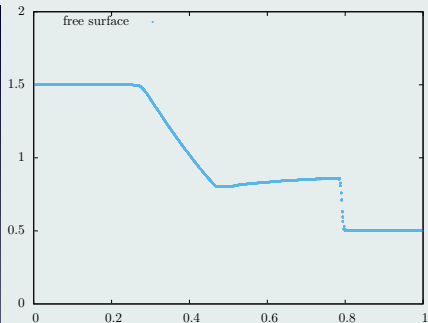


Figure : 3rd-order APLSC-DG on a 5000 cells at $t = 50$: free surface elevation

Dam break problem in cylindrical geometry



(a) Solution map



(b) Solution profile

Figure : 3rd-order APLSC-DG on a 2676 cells mesh at $t = 0.045$: free surface elevation

Rock-wave interaction

Figure : 3rd-order APLSC-DG on a 7000 cells mesh

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A Posteriori Local Subcell Correction (APLSC) technique

- Reformulate DG schemes as subgrid FV-like schemes
- Design an original local subcell correction:
 - preserving the scheme conservation at the subcell scale
 - preserving the very accurate subcell resolution of DG schemes
 - ensuring a maximum or positivity preserving principle at the subcell scale
 - reducing significantly the apparition of spurious oscillations
 - limiting the correction computational effort by not recomputing solution in admissible subcell not lying in the vicinity of a troubled subcell

Applications

- Scalar conservation laws (1D and 2D)
- Euler compressible gas dynamics system (1D and 2D)
- Non-linear shallow water (NSW) system (1D and 2D)
- NSW interactions with a floating object in the arbitrary-Lagrangian-Eulerian (ALE) framework (1D)

Ali Haidar

Ali Haidar

Ongoing work

- Application to 2D total Lagrangian hydrodynamics on curvilinear grids
- Maximum principle DG scheme through subcell reconstructed FCT

Future work

- DoF based adaptive DG scheme through subcell Finite Volume formulation in collaboration with **Raphaël Loubère**
- Application to 2D hydrodynamics and solid dynamics in the ALE framework in collaboration with **Walter Boscheri**
- 2D NSW interactions with a floating object in the ALE framework in collaboration with **Fabien Marche**

Articles on this topic



F. VILAR, *A Posteriori Correction of High-Order DG Scheme through Subcell Finite Volume Formulation and Flux Reconstruction*. JCP, 387:245-279, 2018.



A. HAIDAR, **F. MARCHE** AND **F. VILAR**, *A posteriori Finite-Volume local subcell correction of high-order discontinuous Galerkin schemes for the nonlinear shallow-water equations*. JCP, 452:110902, 2022.



A. HAIDAR, **F. MARCHE** AND **F. VILAR**, *Nonlinear shallow water interactions with a partially immersed object: a robust high-order DG-ALE formulation*. JCP, **Article under revision**.








A. HAIDAR, **F. MARCHE** AND **F. VILAR**, *Numerical approximation of nonlinear shallow-water interacting with a floating object*. **Article in preparation**.



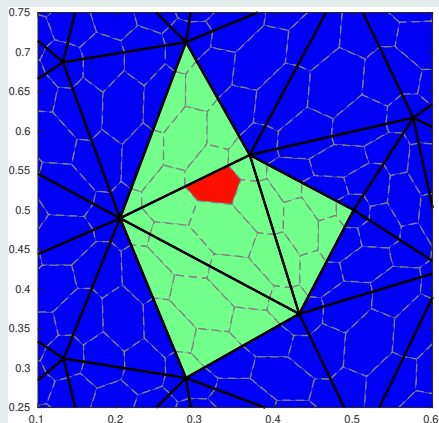
F. VILAR AND **R. ABGRALL**, *A posteriori local subcell correction of DG schemes through Finite Volume reformulation on unstructured grids*. **Article finished, yet to be submitted!!**

Articles on this topic

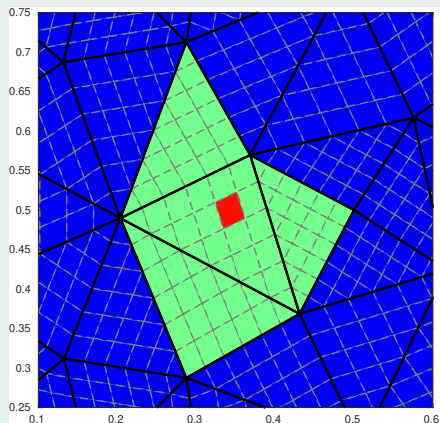
-  **F. VILAR**, *A Posteriori Correction of High-Order DG Scheme through Subcell Finite Volume Formulation and Flux Reconstruction*. JCP, 387:245-279, 2018.
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NAD: neighboring subcells set

linear problems



(a) 4th-order, polygonal subdivision

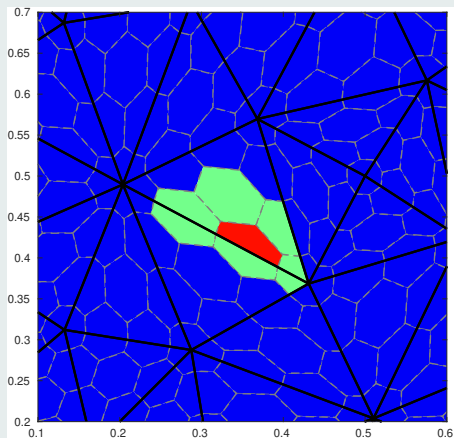


(b) 6th-order, structured subdivision

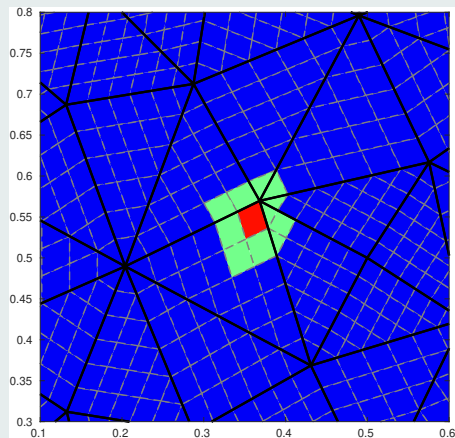
Figure : Neighboring subcells set for the numerical admissibility criterion

NAD: neighboring subcells set

non-linear problems



(a) 4th-order, polygonal subdivision



(b) 6th-order, structured subdivision

Figure : Neighboring subcells set for the numerical admissibility criterion

Cell subdivision: condition number of the projection matrix P_c

	\mathbb{P}^0	\mathbb{P}^1	\mathbb{P}^2	\mathbb{P}^3
Unif. struct. subdiv.	1	4	10.91	31.75
Non-unif. struct. subdiv.	1	4	9.52	29.28
Unif. polyg. subdiv.	1	2.87	8.73	27.89
Non-unif. polyg. subdiv.	1	2.87	8.19	26.94

Table: Projection matrix condition number for different orders and subdivisions